

4. GENERAL SECOND ORDER STRONGLY ELLIPTIC OPERATORS

4.1. Set-Up. Let $\Omega \subseteq \mathbb{R}^n$ be bounded and open. We consider operators of the form

$$(1) \quad A = - \sum_{j,k=1}^n a^{jk}(x) \partial_{x_j} \partial_{x_k} + \sum_j b^j(x) \partial_{x_j} + c(x),$$

where $a^{j,k}$, b^j and c are continuous, real-valued functions, $a^{jk} = a^{kj}$ and

$$(2) \quad c_a |\xi|^2 \leq \sum a^{jk}(x) \xi_j \xi_k \leq C_a |\xi|^2$$

with suitable constants c and C .

We define the forms

$$(3) \quad q_{D/N}^A(u, v) = \int_{\Omega} \sum_{j,k} a^{jk}(x) \partial_{x_k} u(x) \partial_{x_k} \bar{v}(x) + \sum_j b^j(x) u(x) \bar{v}(x) + c(x) u(x) \bar{v}(x) dx$$

with domain $W_0^{1,2}(\Omega) \times W_0^{1,2}(\Omega)$ for q_D^D and $W^{1,2}(\Omega) \times W^{1,2}(\Omega)$ for q_N^D . By Theorem 2.4 they furnish self-adjoint, positive operators $A_{D/N}$.

In a first step we will also assume that all coefficients are constant. Since $a = (a^{jk})$ is real-symmetric, we can write

$$a = s^* d s,$$

where s is an orthogonal matrix and d is diagonal.

We start with the case of a diagonal operator

$$D = - \sum_{k=1}^n d_k^2 \partial_{x_k}^2$$

and denote by $D_{D/N}$ the operators obtained from the corresponding restriction of the form (3).

4.2. Theorem. *Let Ω be a bounded, contented open set in \mathbb{R}^n . Then*

$$N_{D_{D/N}}(\lambda) = \prod_{k=1}^n d_k^{-1} N_{(-\Delta)_{D/N}}(\lambda).$$

Proof. Fix $r > 0$. Following Theorem 3.1 we see that on the interval $I_{rd} = \prod_{k=1}^n [0, rd_k]$ the functions

$$s_j(x) = \prod_{k=1}^n \sqrt{\frac{2}{rd_k}} \sin\left(j_k x_k \frac{\pi}{rd_k}\right), \quad j \in \mathbb{N}^n \text{ and}$$

$$c_j(x) = \prod_{k=1}^n \sqrt{\frac{2}{rd_k}} \cos\left(j_k x_k \frac{\pi}{rd_k}\right), \quad j \in \mathbb{N}_0^n$$

form a complete set of eigenfunctions for the operators D_D and D_N , respectively. We have

$$D_D s_j = \sum_{k=1}^n \frac{j_k^2 \pi^2}{r^2} s_j, \quad j \in \mathbb{N}^n$$

$$D_N c_j = \sum_{k=1}^n \frac{j_k^2 \pi^2}{r^2} c_j, \quad j \in \mathbb{N}_0^n.$$

For the counting functions we therefore obtain:

$$N_{D_{D/N}}(\lambda; I_{rd}) = \left\{ j : \sum_k \frac{j_k^2 \pi^2}{r^2} \leq \lambda \right\} = \left\{ j : \|j\| \leq \frac{r}{\pi} \sqrt{\lambda} \right\} = N_{(-\Delta)_{D/N}}(\lambda; I_r)$$

Here, $N_{(-\Delta)_{D/N}}(\lambda; I_r)$ is the counting function of the Dirichlet/Neumann Laplacian on the cube $I_r = [0, r]^n$, which, as we know, satisfies

$$N_{(-\Delta)_{D/N}}(\lambda; I_{rd}) = \frac{1}{(2\pi)^n} \text{vol}(B(0, 1)) r^n \lambda^{n/2} + O((r^2 \lambda)^{(n-1)/2})$$

For fixed $m \in \mathbb{N}$ we next tile \mathbb{R}^n by the half-open intervals

$$\left[\frac{z_1 d_1}{2^m}, \frac{(z_1 + 1) d_1}{2^m} \right) \times \dots \times \left[\frac{z_n d_n}{2^m}, \frac{(z_n + 1) d_n}{2^m} \right), \quad z = (z_1, \dots, z_n) \in \mathbb{Z}^n.$$

Just as for the Dirichlet/Neumann Laplacian we can then apply the Dirichlet-Neumann bracketing. The difference is that we have replaced the cubes of side length 2^{-m} by intervals of side lengths $2^{-m} d_k$, $k = 1, \dots, n$. Their volume differs by a factor $\prod d_k$ from the cubes' volume. The number of intervals of the form I_d needed to cover Ω from the interior/exterior (in the sense of the sets Ω_m^\mp in the proof of Theorem 3.14) therefore is $\prod d_k^{-1}$ times the number of cubes. This shows the assertion. \square

4.3. Coordinate transforms. Let $T : V \rightarrow U$ be a diffeomorphism of open sets in \mathbb{R}^n and let A be a (differential) operator defined on, say $C_c^\infty(V)$. Then we obtain a (differential) operator $T_* A$ on $C_c^\infty(U)$ by

$$(T_* A u)(x) = A(U \circ T)(T^{-1} x).$$

For later use we compute the first and second order derivatives, assuming for the moment that T is linear, writing $T = (t_{lm})$:

$$\begin{aligned} \partial_{x_k}(u \circ T)(x) &= \partial_{x_k} \left(u \left(\left(\sum_m t_{lm} x_m \right)_l \right) \right) \\ &= \sum_{p=1}^n (\partial_{y_p} u) \left(\left(\sum_m t_{lm} x_m \right)_l \right) t_{pk} \quad \text{and} \\ \partial_{x_j} \partial_{x_k}(u \circ T)(x) &= \partial_{x_j} \left(\sum_{p=1}^n (\partial_{y_p} u) \left(\left(\sum_m t_{lm} x_m \right)_l \right) t_{pk} \right) \\ &= \sum_{p, q=1}^n (\partial_{y_q} \partial_{y_p} u) \left(\left(\sum_m t_{lm} x_m \right)_l \right) t_{qj} t_{pk} \end{aligned}$$

Thus

$$\partial_{x_j} \partial_{x_k}(u \circ T)(T^{-1} x) = \sum_{p, q=1}^n (\partial_{y_q} \partial_{y_p} u)(x) t_{qj} t_{pk}$$

Writing h for the Hessian matrix $(h_{qp})_{qp} = (\partial_{y_q} \partial_{y_p})_{qp}$, we obtain

$$T_*(\partial_{x_j} \partial_{x_k}) = \left(\sum_{q=1}^n \sum_{p=1}^n h_{qp} t_{qj} t_{pk} \right)_{jk} = (T^* h T)_{jk}.$$

4.4. Example. We consider the principal part A^H of the operator A in 4.1 with constant coefficients:

$$A^H = - \sum_{jk} a^{jk} \partial_{x_j} \partial_{x_k}.$$

As in 4.1(3) we assume that $a = s^* d s$ for an orthogonal matrix s and a diagonal matrix $d = \text{diag}(d_1^2, \dots, d_n^2)$. We also suppose that the boundary of Ω is C^2 .

We derive from (4.3) that

$$\begin{aligned}
s_*A^H &= -\sum_{jk} a^{jk}(s^*hs)_{jk} = -\sum_{jk} a^{jk}(s^*hs)_{jk} \\
&= -\operatorname{Tr}(as^*hs) = -\operatorname{Tr}(sas^*h) = \operatorname{Tr}(dh) \\
&= -\sum_k d_k^2 \partial_{x_k}^2 = D
\end{aligned}$$

with the operator D defined in 4.1.

By $A_{D/N}^H$ we denote the self-adjoint, positive operators obtained from the forms

$$q_{D/N}(u, v) = \sum_{jk} \int_{\Omega} a^{jk} \partial_{x_j} u \partial_{x_k} \bar{v} \, dx$$

with domains $W_0^{1,2}(\Omega) \times W_0^{1,2}(\Omega)$ and $W^{1,2}(\Omega) \times W^{1,2}(\Omega)$, respectively. These operators have compact resolvent and therefore their spectrum consists of eigenvalues $0 \leq \lambda_k \rightarrow +\infty$. Since the boundary is C^2 , the operators A_D^H and A_N^H act on their domains as the operator A^H .

Now we make the following observation: Suppose v is an eigenfunction for the eigenvalue λ of A^H . Then $u = v \circ s^{-1}$ is an eigenfunction for the eigenvalue λ of s_*A^H :

$$(s_*A^H)u(x) = A^H(u \circ s)(s^{-1}(x)) = (A^H v)(s^{-1}(x)) = \lambda v(s^{-1}(x)) = \lambda u(x).$$

Conversely, if u is an eigenfunction for the eigenvalue λ of s_*A^H , then $v = u \circ s$ is an eigenfunction for A^H . Hence the spectra and also the counting function for A^H and s_*A^H coincide. We therefore obtain

$$N_{A_{D/N}^H}(\lambda) = N_{s_*A_{D/N}^H}(\lambda) = N_{D_{D/N}}(\lambda) = \frac{1}{\prod d_k} N_{-\Delta_{D/N}}(\lambda) = \frac{1}{\sqrt{\det a}} N_{-\Delta_{D/N}}(\lambda).$$

4.5. The constant coefficient case. Assume Weyl asymptotics hold for the Dirichlet and the Neumann Laplacian on Ω . For the Dirichlet and Neumann realizations $A_{D/N}$ of the operator A in 4.1 with constant coefficients we have.

$$N_{A_{D/N}}(\lambda) = (\det a)^{-1/2} \frac{\operatorname{vol} \Omega \operatorname{vol} B(0, 1)}{(2\pi)^n} \lambda^{n/2} + o(\lambda^{n/2}).$$

Proof. Consider the forms

$$q_{D/N}(u, v) = \int_{\Omega} \sum_{jk} a^{jk} \partial_{x_j} u \partial_{x_k} \bar{v} + \sum_{j=1}^n b^j \partial_{x_j} u \bar{v} + c u \bar{v} \, dx$$

with domains $W_0^{1,2}(\Omega) \times W_0^{1,2}(\Omega)$ and $W^{1,2}(\Omega) \times W^{1,2}(\Omega)$, respectively. These forms are semi-bounded: For every j and every $\varepsilon > 0$ there exists a C_ε such that

$$b^j \partial_{x_j} u \bar{v} \leq |b^j| |\partial_{x_j} u \bar{v}| \leq |b^j| \varepsilon |\partial_{x_j} u|^2 + C_\varepsilon |v|^2.$$

By assumption $\sum_{j,k} a^{jk} \xi_j \xi_k \geq c_a |\xi|^2$ for some constant $c_a > 0$. We deduce that, for every $\varepsilon > 0$, we find a C_ε , depending on ε , b and c with

$$q_{D/N}(u, u) \geq (c_a - \varepsilon) \|\nabla u\|_{L^2(\Omega)}^2 - C_\varepsilon \|u\|_{L^2(\Omega)}^2.$$

It is also closed, so that we obtain the self-adjoint operators $A_{D/N}$ from Theorem 2.4. They have compact resolvents and correspondingly, a discrete set of real eigenvalues tending to $+\infty$.

In order to determine the asymptotic of the respective counting functions of A from those of A^H , we compare the forms. The estimate

$$\left| \sum_{j=1}^n \int b^j \partial_{x_j} u \bar{u} + cu \bar{u} dx \right| \leq \varepsilon \|\nabla u\|_{L^2(\Omega)} + C_\varepsilon \|u\|^2$$

for arbitrary $\varepsilon > 0$ and suitable C_ε implies that, also for arbitrary $\varepsilon > 0$ and suitable C_ε ,

$$\begin{aligned} q_{D/N}(u, u) &\leq (1 + \varepsilon) q_{D/N}^{A^H}(u, u) + C_\varepsilon \|u\|_{L^2(\Omega)} \quad \text{and} \\ q_{D/N}(u, u) &\geq (1 - \varepsilon) q_{D/N}^{A^H}(u, u) - C_\varepsilon \|u\|_{L^2(\Omega)}. \end{aligned}$$

Since the forms $q_{D/N}$ and $q_{D/N}^{A^H}$ have the same domains, we infer from the minimax principle that

$$(1) \quad \lambda_k^{A_{D/N}} \leq (1 + \varepsilon) \lambda_k^{A_{D/N}^{A^H}} + C_\varepsilon \quad \text{and}$$

$$(2) \quad \lambda_k^{A_{D/N}} \geq (1 - \varepsilon) \lambda_k^{A_{D/N}^{A^H}} - C_\varepsilon.$$

Hence $\lambda_k^{A_{D/N}^{A^H}} \leq \lambda$ implies that $(1 - \varepsilon) \lambda_k^{A_{D/N}^{A^H}} - C_\varepsilon \leq \lambda$ which in turn shows that $\lambda_k^{A_{D/N}^{A^H}} \leq \frac{1}{1 - \varepsilon} \lambda + \frac{C_\varepsilon}{1 - \varepsilon}$. We conclude from (2) that

$$\begin{aligned} N_{A_{D/N}}(\lambda) &\leq N_{A_{D/N}}\left(\frac{1}{1 - \varepsilon} \lambda + \frac{C_\varepsilon}{1 - \varepsilon}\right) \\ &= (\det a)^{-1/2} N_{-\Delta_{D/N}}\left(\frac{1}{1 - \varepsilon} \lambda + \frac{C_\varepsilon}{1 - \varepsilon}\right) + O\left(\left(\frac{1}{1 - \varepsilon} \lambda + \frac{C_\varepsilon}{1 - \varepsilon}\right)^{\frac{n-1}{2}}\right) \end{aligned}$$

Similarly, (1) implies that

$$\begin{aligned} N_{A_{D/N}}(\lambda) &\geq N_{A_{D/N}}\left(\frac{1}{1 + \varepsilon} \lambda - \frac{C_\varepsilon}{1 + \varepsilon}\right) \\ &= (\det a)^{-1/2} N_{-\Delta_{D/N}}\left(\frac{1}{1 + \varepsilon} \lambda - \frac{C_\varepsilon}{1 + \varepsilon}\right) + O\left(\left(\frac{1}{1 + \varepsilon} \lambda - \frac{C_\varepsilon}{1 + \varepsilon}\right)^{\frac{n-1}{2}}\right) \end{aligned}$$

This implies the desired estimate: In fact, we know that $C_\varepsilon \leq C/\varepsilon$ for suitable C . Let $f : \mathbb{R}_{>0} \rightarrow \mathbb{R}_{>0}$ be a function satisfying

$$f(\lambda) \geq (1 - \varepsilon) \lambda^{n/2} - \frac{C}{\varepsilon} o(\lambda^{n/2}) \quad \text{and} \quad f(\lambda) \leq (1 + \varepsilon) \lambda^{n/2} + \frac{C}{\varepsilon} o(\lambda^{n/2})$$

Suppose $\liminf_{\lambda \rightarrow \infty} f(\lambda) \lambda^{-n/2} - 1 < -\delta$ for some $\delta > 0$. Choosing $\varepsilon = \delta/2$ this would imply that for some sequence $\lambda^{(k)}$ we would have, for all k ,

$$-\frac{\delta}{2} + \frac{2C}{\delta} (\lambda^{(k)})^{-1/2} < -\delta,$$

which is not possible. In the same way we can treat $\limsup_{\lambda \rightarrow \infty} f(\lambda) \lambda^{-n/2} - 1$. \square

4.6. Variable coefficients, Dirichlet boundary conditions. Let Ω be open and contented. For the Dirichlet realization A_D of the operator A in 4.1 with variable coefficients we then have.

$$N_{A_D}(\lambda) = \frac{\text{vol}_g \Omega \text{vol} B(0, 1)}{(2\pi)^n} \lambda^{n/2} + o(\lambda^{n/2}),$$

where $\text{vol}_g \Omega$ is the volume of Ω measured in the Riemannian metric $g = a^{-1}$.

Proof. For fixed $m \in \mathbb{N}$ we next tile \mathbb{R}^n by the half-open intervals

$$I_z = \left[\frac{z_1 d_1}{2^m}, \frac{(z_1 + 1) d_1}{2^m} \right) \times \left[\frac{z_n d_n}{2^m}, \frac{(z_n + 1) d_n}{2^m} \right), \quad z = (z_1, \dots, z_n) \in \mathbb{Z}^n.$$

On each interval I_z consider the form

$$q_D(u, v) = \int_{I_z} \sum_{jk} a^{jk} \partial_{x_j} u \partial_{x_k} \bar{v} + \sum_{j=1}^n b^j \partial_{x_j} u \bar{v} + c u \bar{v} dx$$

with domain $W_0^{1,2}(I_z) \times W_0^{1,2}(I_z)$. This form is closed and semi-bounded by the same arguments as in the proof of 4.5. It defines a selfadjoint operator A_D with discrete real spectrum consisting of eigenvalues $\lambda_k^{A_D} j$ tending to ∞ .

Denote by x_z the midpoint of I_z . We can then compare the forms q_D with the constant coefficient forms

$$q_D^z(u, v) = \int_{I_z} \sum_{jk} a^{jk}(x_z) \partial_{x_j} u \partial_{x_k} \bar{v} + \sum_{j=1}^n b^j(x_z) \partial_{x_j} u \bar{v} + c(x_z) u \bar{v} dx$$

Let $\varepsilon > 0$ be given. By taking m large (and hence the diameter of the intervals small) we can achieve that, with suitable $C_\varepsilon > 0$,

$$\begin{aligned} q_D(u, u) &\leq (1 + \varepsilon) q_D^z(u, u) + C_\varepsilon \|u\|_{L^2(\Omega)} \quad \text{and} \\ q_D(u, u) &\geq (1 - \varepsilon) q_D^z(u, u) - C_\varepsilon \|u\|_{L^2(\Omega)}. \end{aligned}$$

As in the proof of 4.5 we conclude that the counting functions on each I_z satisfy

$$\begin{aligned} N_{A_D}(\lambda) &\geq (1 - \varepsilon) (\det a_z)^{-1/2} (2\pi)^{-n} \text{vol } I_z \text{vol } B(0, 1) \lambda^{n/2} - C_\varepsilon \cdot o(\lambda^{n/2}) \quad \text{and} \\ N_{A_D}(\lambda) &\leq (1 + \varepsilon) (\det a_z)^{-1/2} (2\pi)^{-n} \text{vol } I_z \text{vol } B(0, 1) \lambda^{n/2} + C_\varepsilon \cdot o(\lambda^{n/2}) \end{aligned}$$

where $a_c = (a^{jk}(x_z))_{jk}$. Dirichlet-Neumann bracketing and taking the limit as $m \rightarrow \infty$ then implies that

$$N_{A_D}(\lambda) = \int_{\Omega} (\det a(x))^{-1/2} dx \frac{\text{vol } B(0, 1)}{(2\pi)^n} \lambda^{n/2} + o(\lambda^{n/2}).$$

Finally we note that

$$\int_{\Omega} (\det a(x))^{-1/2} dx = \int_{\Omega} \sqrt{\det a(x)^{-1}} dx = \text{vol}_g(\Omega)$$

for the Riemannian metric $g = a^{-1}$. □

4.7. Remark. Second order operators on manifolds. Let Ω be a smooth manifold of dimension n , and let A be a strongly elliptic second order operator that locally is of the form in 4.1.

According to Whitehead [12] every smooth manifold admits a triangulation. For each simplicial set we can consider the corresponding operator on a simplex in \mathbb{R}^n . On the simplex, we can derive Weyl asymptotics as in 4.6 for the local operator.

I expect that, via Dirichlet-Neumann bracketing we then obtain the Weyl asymptotics on the manifold.