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# DECOUPLING FOR SCHATTEN CLASS OPERATORS IN THE SETTING OF QUANTUM HARMONIC ANALYSIS

Workshop on Quantum Harmonic Analysis - Hannover Helge Jørgen Samuelsen August 06, 2024 **Notation and Background** 

**Classical results** 

**Quantum Harmonic Analysis** 

Main results

Sketch of the proof



# Notation and Background - Notation

Let  $\mathscr{S}(\mathbb{R}^{2d})$  denote the Schwartz class of rapidly decaying smooth functions, and  $\mathscr{S}'(\mathbb{R}^{2d})$  its dual space of tempered distributions.

#### Symplectic Fourier transform

For  $\psi \in \mathscr{S}(\mathbb{R}^{2d})$  we denote the symplectic Fourier transform of  $\psi$  at a point  $\zeta \in \mathbb{R}^{2d}$  by

$$\mathcal{F}_{\sigma}(\psi)(\zeta) = \int_{\mathbb{R}^{2d}} e^{-2\pi i \sigma(\zeta,z)} \psi(z) \, dz,$$

where  $\sigma$  is the symplectic form on  $\mathbb{R}^{2d}$ . For  $\tau \in \mathscr{S}'(\mathbb{R}^{2d})$ , we extend the sympletic Fourier transform to  $\mathscr{S}'(\mathbb{R}^{2d})$  through,

$$\langle \mathcal{F}_{\sigma}(\tau),\psi\rangle_{\mathscr{S}',\mathscr{S}}=\langle \tau,\mathcal{F}_{\sigma}(\psi)\rangle_{\mathscr{S}',\mathscr{S}},\quad\forall\psi\in\mathscr{S}(\mathbb{R}^{2d}),$$

where a sesquilinear dual pairing is used.

# Notation and Background - Classical Decoupling

Let us introduces the classical decoupling constant.

#### Definition

Let  $\Omega \subset \mathbb{R}^{2d}$ , and  $\mathcal{P}_{\Omega}$  be a partition of  $\Omega$ . We define the classical decoupling constant  $\mathcal{D}_{p,q}^{\mathscr{C}}(\mathcal{P}_{\Omega})$  to be the smallest constant for which the inequality

$$\left\|\sum_{\theta\in\mathcal{P}_\Omega}f_\theta\right\|_{L^p(\mathbb{R}^{2d})}\leq\mathcal{D}^{\mathscr{C}}_{p,q}(\mathcal{P}_\Omega)\left(\sum_{\theta\in\mathcal{P}_\Omega}\|f_\theta\|^q_{L^p(\mathbb{R}^{2d})}\right)^{\frac{1}{q}},$$

holds for any collection of functions with supp  $\mathcal{F}_{\sigma}(f_{\theta}) \subseteq \theta$ .

The goal of decoupling is to achieve asymptotic bounds for the decoupling constant when considering finer and finer partitions, i.e.  $\#\mathcal{P}_{\Omega} \to \infty$ . It is mostly q = 2 which is considered in the literature.

# Notation and Background - Basic estimates

# **Square function estimates**

By Minkowski's inequality for p > q, it follows that

$$\left\| \left( \sum_{\theta \in \mathcal{P}_{\Omega}} |f_{\theta}|^{q} \right)^{\frac{1}{q}} \right\|_{L^{p}(\mathbb{R}^{2d})} \leq \left( \sum_{\theta \in \mathcal{P}_{\Omega}} \|f_{\theta}\|_{L^{p}(\mathbb{R}^{2d})}^{q} \right)^{\frac{1}{q}}.$$

In particular, for p > 2, every square function estimate gives rise to a decoupling result for q = 2. We may therefore think of decoupling as a weaker square function estimate.

# Trivial decoupling constant for finite partitions

Assume the partition  $\mathcal{P}_{\Omega}$  consist of *N* elements. Then Hölder's inequality yields

$$\left\|\sum_{\theta\in\mathcal{P}_{\Omega}}f_{\theta}\right\|_{L^{p}(\mathbb{R}^{2d})}\leq\sum_{\theta\in\mathcal{P}_{\Omega}}\|f_{\theta}\|_{L^{p}(\mathbb{R}^{2d})}\leq N^{\frac{1}{q'}}\left(\sum_{\theta\in\mathcal{P}_{\Omega}}\|f_{\theta}\|_{L^{p}(\mathbb{R}^{2d})}^{q}\right)^{\frac{1}{q}}$$



# Classical results - Title of frame

Consider the truncated paraboloid

$$\mathbb{P}^{2d-1} := \{ (\xi, |\xi|^2) \in \mathbb{R}^{2d} : \xi \in \mathbb{R}^{2d-1}, |\xi_i| \le 1 \text{ for } 1 \le i \le 2d-1 \},$$

and cover the  $\delta\text{-neighbourhood}$  by

$$heta=\{(\xi,|\xi|^2+t):\xi\in C_ heta,|t|<\delta\},$$

where  $\{C_{\theta}\}$  is a partition of  $[-1, 1]^{2d-1}$  by cubes of side length  $\delta^{\frac{1}{2}}$ .

#### Theorem (Bourgain-Demeter '15)

For  $0 < \delta < 1$  and any  $\varepsilon > 0$  there exists a constant  $C_{\varepsilon} > 0$  such that if  $2 \le p \le (4d + 2)/(2d - 1)$ , then

$$\mathcal{D}^{\mathscr{C}}_{p,2}(\mathcal{P}_{\mathcal{N}_{\mathbb{P}^{2d-1}}(\delta)}) \leq C_{\varepsilon}\delta^{-\varepsilon},$$

*while if*  $p \ge (4d + 2)/(2d - 1)$ *, then* 

$$\mathcal{D}_{p,2}^{\mathscr{C}}(\mathcal{P}_{\mathcal{N}_{\mathbb{P}^{2d-1}}(\delta)}) \leq C_{\varepsilon}\delta^{-\varepsilon+rac{2d+1}{2p}-rac{2d-1}{4}}.$$

The proved the result for compact hypersurfaces with non-vanishing Gaussian curvature.



# **Classical results -** Figure of the $\delta$ -neighbourhood



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#### **Classical results** - Decoupling for the moment curve

Consider the moment curve  $\Gamma:[0,1] \to \mathbb{R}^{2d}$  given by

$$\Gamma(t)=(t,t^2,\ldots,t^{2d-1},t^{2d}).$$

For each  $0 < \delta < 1$ , let  $\mathcal{P}(\delta)$  be a dyadic partition

$$[0,1] = igcup_{J\in \mathcal{P}(\delta)} J$$

where each interval  $|J| = 2^{\lceil \log_2(\delta) \rceil}$ . Moreover, for each interval J let  $c_J$  be the centre of J and denote by  $\mathcal{U}_J$  the parallelepiped centred at  $\Gamma(c_J)$  of dimensions  $|J| \times |J|^2 \times \ldots \times |J|^{2d}$  where side number j is parallel to  $\Gamma^{(j)}(c_J)$ .

#### Theorem (Bourgain-Demeter-Guth '16)

For each  $d \in \mathbb{N}$  and every  $\varepsilon > 0$ , there exists a finite constant  $C_{d,\varepsilon} > 0$  such that

 $\mathcal{D}^{\mathscr{C}}_{2d(2d+1),2}(\mathcal{P}(\delta)) \leq C_{d,\varepsilon}\delta^{-\varepsilon},$ 

for every  $0 < \delta < 1$ .

# **Classical results - Application to exponential sums**

# Lemma (Reverse Hölder)

Let  $|\mathcal{P}_{\Omega}| = N \in \mathbb{N}$ . Then for each  $1 \leq p, q \leq \infty$  and  $a \in \ell^q(\{1, \dots, N\})$  and  $R \gtrsim \max |\theta|^{-1}$ ,

$$\left( \int_{B_R(0)} \left| \sum_{\theta \in \mathcal{P}_\Omega} a_\theta e^{2\pi i \sigma(x,\xi_\theta)} \right|^p dx \right)^{\frac{1}{p}} \lesssim \mathcal{D}_{p,q}^{\mathscr{C}}(\mathcal{P}_\Omega) \|a\|_{\ell^q},$$

where  $\xi_{\theta} \in \theta$ .

#### **Corollary (Application to the moment curve)**

For each  $1 \le i \le N$ , let  $t_i \in ((i-1)/N, i/N]$ . For each  $R \ge N^{2d}$ , and every  $a \in \ell^2(\{1, \ldots, N\})$  and every  $p \ge 2$  there exists  $C_{\varepsilon} > 0$  such that

$$\left(\int_{B_R(0)}\left|\sum_{l=1}^N a_l e^{2\pi i \sigma(x,\Gamma(t_l))}\right|^p dx\right)^{\frac{1}{p}} \leq C_{\varepsilon} \left(N^{\varepsilon} + N^{\frac{1}{2} - \frac{2d(2d+1)}{2p} + \varepsilon}\right) \|a\|_{\ell^2}$$



# **Classical results - Sketch of proof for reverse Hölder**

Proof.

Let  $\psi \in \mathscr{S}(\mathbb{R}^{2d})$  be such that  $|\psi| \ge 1$  on  $B_1(0)$  and  $\mathcal{F}_{\sigma}(\psi)$  is supported on  $B_1(0)$ . Then the function

$${\it F}_{ heta}(z) = {\it a}_{ heta} \psi_{\it R}(z) e^{2\pi i \sigma(z, \xi_{ heta})}$$

has its symplectic Fourier transform support on  $\theta$ . Here  $\psi_R(z) = \psi(z/R)$ . Decoupling then gives

$$\left(\int_{B_{R}(0)}\left|\sum_{\theta\in\mathcal{P}_{\Omega}}a_{\theta}e^{2\pi i\sigma(z,\xi_{\theta})}\right|^{p}dz\right)^{\frac{1}{p}} \leq \left\|\sum_{\theta\in\mathcal{P}_{\Omega}}F_{\theta}\right\|_{L^{p}} \leq \mathcal{D}_{p,q}^{\mathscr{C}}(\mathcal{P}_{\Omega})\left(\sum_{\theta\in\mathcal{P}_{\Omega}}\|a_{\theta}\psi_{R}\|_{L^{p}}^{q}\right)^{\frac{1}{q}} = R^{\frac{2d}{p}}\|\psi\|_{L^{p}}\mathcal{D}_{p,q}^{\mathscr{C}}(\mathcal{P}_{\Omega})\|a\|_{\ell^{q}}.$$

The proof is concluded by recalling that  $|B_r(0)| \approx R^{2d}$ .

# **Time-frequency Analysis**

Let  $\rho$  denote the symmetric time-frequency shift given by

$$\rho(x,\xi)f(t) = e^{-\pi i x \cdot \xi} e^{2\pi i \xi \cdot t} f(t-x), \quad (x,\xi) \in \mathbb{R}^{2d}, \ f \in L^2(\mathbb{R}^d)$$

Associated to  $\rho$  is the cross-ambiguity function. For  $f, g \in L^2(\mathbb{R}^d)$ , the cross-ambiguity function of f and g is defined at a point  $z \in \mathbb{R}^{2d}$  by

$$egin{aligned} \mathcal{A}(f,g)(z) &= \langle f,
ho(z)g
angle_{L^2}\in L^2(\mathbb{R}^{2d}). \ & \mathcal{W}(f,g)(z) &= \mathcal{F}_\sigma\left(\mathcal{A}(f,g)
ight)(z). \end{aligned}$$

# **Quantum Harmonic Analysis - Operators**

#### **Schatten Class**

Let  $\mathcal{K}$  denote the compact operators on  $L^2(\mathbb{R}^d)$ . By the singular value decomposition, for each  $\mathcal{T} \in \mathcal{K}$  we can write it as \_\_\_\_\_

$$\mathcal{T} = \sum_{n \in \mathbb{N}} s_n(\mathcal{T}) \varphi_n \otimes \psi_n,$$

where  $\{s_n(T)\}\)$  are the singular values of T, and  $\{\varphi\}$ ,  $\{\psi_n\}\)$  are two orthonomal families in  $L^2$ . We denote the Schatten *p*-class of compact operators on  $L^2(\mathbb{R}^d)$  by

$$\mathcal{S}^{p} = \{ T \in \mathcal{K} : \{ s_{n}(T) \} \in \ell^{p} \}$$

and equip it with the norm

$$\|T\|_{\mathcal{S}^p} = \left(\sum_{n\in\mathbb{N}} |s_n(T)|^p\right)^{rac{1}{p}},$$

which turn it into a Banach space.



# **Quantum Harmonic Analysis** - Operators

# Weyl quantization

For each  $\tau \in \mathscr{S}'(\mathbb{R}^{2d})$ , we can associate it to a continuous linear operator from  $\mathscr{S}(\mathbb{R}^d)$  to  $\mathscr{S}'(\mathbb{R}^d)$ , denoted  $L_{\tau}$ , by the cross-Wigner distribution,

$$\langle L_{\tau}f,g\rangle_{\mathscr{S}',\mathscr{S}}=\langle \tau,\mathcal{W}(g,f)\rangle_{\mathscr{S}',\mathscr{S}},$$

for all  $f, g \in \mathscr{S}(\mathbb{R}^d)$ . Here we use a sesquilinear dual pairing. This is known as the Weyl quantization of  $\tau$ , and  $\tau$  is called the Weyl symbol of  $L_{\tau}$ . In fact, every continuous linear operator from  $\mathscr{S}(\mathbb{R}^d)$  to  $\mathscr{S}'(\mathbb{R}^d)$  has a Weyl symbol.

#### Schwartz and tempered distribution classes of operators

Define the Schwartz class of operators by

$$\mathfrak{S} = \{L_{\varphi} : \varphi \in \mathscr{S}(\mathbb{R}^{2d})\} \subseteq \mathcal{S}^1,$$

and its dual space

$$\mathfrak{S}'\cong\mathcal{L}\left(\mathscr{S}(\mathbb{R}^d),\mathscr{S}'(\mathbb{R}^d)
ight).$$



# **Quantum Harmonic Analysis** - Convolutions

# Convolutions

We then define an operator-operator convolution by

 $T \star S(z) := \operatorname{tr}(T\alpha_z(PSP)),$ 

and a function-operator convolution by

$$f \star T = T \star f := \int_{\mathbb{R}^{2d}} f(z) \alpha_z(T) \, dz,$$

where the integral is considered weakly. Here  $\alpha_z(T) = \rho(z)T\rho(-z)$  and Pf(t) = f(-t) is the Parity operator.

#### **Theorem (Werner-Young)**

Let  $1 \leq p, q, r \leq \infty$  be such that  $1 + r^{-1} = p^{-1} + q^{-1}$ . If  $f \in L^p(\mathbb{R}^{2d})$ ,  $T \in S^p$  and  $S \in S^q$ , then  $f \star S \in S^r$  and  $T \star S \in L^r(\mathbb{R}^{2d})$ . Moreover, there are the norm bounds

 $\begin{aligned} \|f \star S\|_{\mathcal{S}^r} \leq \|f\|_{L^p} \|S\|_{\mathcal{S}^q}, \\ \|T \star S\|_{L^r} \leq \|T\|_{\mathcal{S}^p} \|S\|_{\mathcal{S}^q}. \end{aligned}$ 



# Quantum Harmonic Analysis - The basics of Fourier-Wigner transform

#### **Fourier-Wigner transform**

For  $T \in S^1$ , and  $z \in \mathbb{R}^{2d}$ , we define the Fourier-Wigner (or Fourier-Weyl) transform as

 $\mathcal{F}_W(\mathcal{T})(z) = \operatorname{tr}(\mathcal{T}\rho(-z)) \in C_0(\mathbb{R}^{2d}).$ 

It also extends to an isomorphism from  $S^2$  to  $L^2(\mathbb{R}^{2d})$  due to a result of Pool. Since  $\mathfrak{S} \subseteq S^1$ , the Fourier-Wigner is well-defined and can be extended by duality to  $\mathfrak{S}'$ .

#### Theorem (Keyl-Kiukas-Werner '16)

Let  $S \in \mathfrak{S}$ ,  $\psi \in \mathscr{S}(\mathbb{R}^{2d})$ ,  $\tau \in \mathscr{S}'(\mathbb{R}^{2d})$  and  $A \in \mathfrak{S}'$ . Then relations

$\mathcal{F}_{\sigma}(S \star A) = \mathcal{F}_{W}(S)\mathcal{F}_{W}(A),$	$\mathcal{F}_{W}(\psi\star A)=\mathcal{F}_{\sigma}(\psi)\mathcal{F}_{W}(A),$
$\mathcal{F}_{\sigma}(\psi st  au) = \mathcal{F}_{\sigma}(\psi) \mathcal{F}_{\sigma}( au),$	$\mathcal{F}_W(\boldsymbol{S}\star  au) = \mathcal{F}_W(\boldsymbol{S})\mathcal{F}_\sigma( au),$

hold.



# Quantum Harmonic Analysis - Quantum decoupling

#### **Extending decoupling to operators**

We can now extend decoupling to operators with the following definition.

#### Definition

Let  $\Omega \subset \mathbb{R}^{2d}$ , and  $\mathcal{P}_{\Omega}$  be a partition of  $\Omega$ . We define the quantum decoupling constant  $\mathcal{D}_{p,q}^{\mathscr{Q}}(\mathcal{P}_{\Omega})$  to be the smallest constant for which the inequality

$$\left|\sum_{ heta\in\mathcal{P}_\Omega} extsf{T}_ heta
ight|_{\mathcal{S}^p} \leq \mathcal{D}^{\mathscr{Q}}_{p,q}(\mathcal{P}_\Omega) \left(\sum_{ heta\in\mathcal{P}_\Omega} \| extsf{T}_ heta\|^q_{\mathcal{S}^p}
ight)^rac{1}{q},$$

holds for any collection of operators  $\{T_{\theta}\}_{\theta \in \mathcal{P}_{\Omega}} \subset \mathcal{L}(\mathscr{S}(\mathbb{R}^{d}); \mathscr{S}'(\mathbb{R}^{d}))$  with  $\operatorname{supp}(\mathcal{F}_{W}(T_{\theta})) \subseteq \theta$ .

# Main results - Equivalence of decoupling constants

#### Equivalence of classical and quantum decoupling constants

We are able to show the following equivalence of decoupling under the assumption of  $\Omega$  being a bounded set.

#### Theorem

Let  $\Omega \subset \mathbb{R}^{2d}$  be a bounded set, and let  $\mathcal{P}_{\Omega}$  be a partition of  $\Omega$ . Then there exists  $C = C(\Omega) \ge 1$  such that

$$C^{-1}\mathcal{D}^{\mathscr{C}}_{p,q}(\mathcal{P}_\Omega) \leq \mathcal{D}^{\mathscr{Q}}_{p,q}(\mathcal{P}_\Omega) \leq C\,\mathcal{D}^{\mathscr{C}}_{p,q}(\mathcal{P}_\Omega)$$

holds for any  $1 \leq p, q \leq \infty$ .

#### Remark

The constant C in the theorem is independent of p, q and the partition  $\mathcal{P}_{\Omega}$ .

#### Main results - Quantum decoupling for the Paraboloid

Consider the truncated paraboloid

$$\mathbb{P}^{2d-1} := \{ (\xi, |\xi|^2) \in \mathbb{R}^{2d} : \xi \in \mathbb{R}^{2d-1}, |\xi_i| \le 1 \text{ for } 1 \le i \le 2d-1 \},$$

and again cover the  $\delta$ -neighbourhood of  $\mathbb{P}^{2d-1}$  by

$$heta=\{(\xi,|\xi|^2+t):\xi\in \mathcal{C}_ heta,|t|<\delta\},$$

where  $\{C_{\theta}\}$  is a partition of  $[-1, 1]^{2d-1}$  by cubes of side length  $\delta^{\frac{1}{2}}$ .

# $\ell^2 \mathcal{S}^p$ -quantum decoupling

For  $0 < \delta < 1$  and any  $\varepsilon > 0$  there exists a constant  $C_{\varepsilon} > 0$  such that if  $2 \le p \le (4d + 2)/(2d - 1)$ , then

$$\mathcal{D}^{\mathscr{Q}}_{p,2}(\mathcal{P}_{\mathcal{N}_{\mathbb{P}^{2d-1}}(\delta)}) \leq C_{\varepsilon}\delta^{-\varepsilon},$$

while if  $p \ge (4d+2)/(2d-1)$ , then

$$\mathcal{D}^{\mathscr{Q}}_{p,2}(\mathcal{P}_{\mathcal{N}_{\mathbb{P}^{2d-1}}(\delta)}) \leq C_{\varepsilon}\delta^{-\varepsilon + \frac{2d+1}{2p} - \frac{2d-1}{4}}$$



# Main results - A convolution result

# **Two Convolution results**

# Proposition

Let  $\Omega$  be a bounded set in  $\mathbb{R}^{2d}$ . There exist  $L^2$ -normalised  $g, h \in \mathscr{S}(\mathbb{R}^d)$  and  $B \in \mathfrak{S}$  such that:

- i) If  $\tau \in \mathscr{S}'(\mathbb{R}^{2d})$  and  $\mathcal{F}_{\sigma}(\tau)$  is supported on  $\Omega$ , then  $\tau = \tau \star (g \otimes h) \star B$ .
- ii) If  $T \in \mathfrak{S}'$  and  $\mathcal{F}_W(T)$  is supported on  $\Omega$ , then  $T = T \star (g \otimes h) \star B$ .

# Corollary

Let  $\Omega$  be a bounded set, and  $1 \le p \le \infty$ . There exists  $L^2$ -normalised  $g, h \in \mathscr{S}(\mathbb{R}^d)$  and  $C = C(\Omega) > 0$  such that

i) If  $\tau \in \mathscr{S}'(\mathbb{R}^{2d})$  and  $\mathcal{F}_{\sigma}(\tau)$  is supported on  $\Omega$ , then

 $\| au\|_{L^p(\mathbb{R}^{2d})} \leq C(\Omega) \| au \star (g \otimes h)\|_{\mathcal{S}^p}.$ 

ii) If  $T \in \mathfrak{S}'$  and  $\mathcal{F}_W(T)$  is supported on  $\Omega$ , then

$$\|T\|_{\mathcal{S}^p} \leq C(\Omega) \|T\star (g\otimes h)\|_{L^p(\mathbb{R}^{2d})}.$$



# Sketch of the proof - Proof of Convolution result

- ▶ There exist  $z_0 \in \Omega$  and R > 0 such that  $\Omega \subseteq B_R(z_0)$ . Let  $\Psi \in C_c^{\infty}(\mathbb{R}^{2d})$  be a smooth cut-off function such that  $\Psi \equiv 1$  on  $B_R(z_0)$  and supported on  $B_{2R}(z_0)$ . Moreover, let  $g, h \in \mathscr{S}(\mathbb{R}^d)$  be such that  $|\mathcal{A}(g, h)(z)| \ge \delta > 0$  on  $B_{2R}(z_0)$ .
- Then, for each  $z \in \Omega$ ,

$$\Psi(z)=rac{\mathcal{A}(g,h)(z)}{\mathcal{A}(g,h)(z)}\Psi(z)=\mathcal{F}_W(g\otimes h)(z)\mathcal{F}_W(B)(z),$$

where the operator *B* is given by

$$B = \mathcal{F}_W^{-1}\left(\frac{\Psi}{\mathcal{A}(g,h)}\right) = \int_{\mathbb{R}^{2d}} \frac{1}{\mathcal{A}(g,h)(z)} \Psi(z) \rho(z) \, dz.$$

For each  $T \in \mathfrak{S}'$  such that supp  $\mathcal{F}_W(T) \subseteq \Omega$ , it follows that

$$\mathcal{F}_W(T) = \mathcal{F}_W(T)\Psi = \mathcal{F}_W(T)\mathcal{F}_W(g\otimes h)\mathcal{F}_W(B) = \mathcal{F}_W(T\star(g\otimes h)\star B),$$

and the result follows as  $\mathcal{F}_W$  is an isomorphism from  $\mathfrak{S}'$  to  $\mathscr{S}'(\mathbb{R}^{2d})$ .

# Sketch of the proof - Proof of decoupling equivalence

Let {*T*<sub>θ</sub>} be a collection for which supp  $\mathcal{F}_W(T_\theta) \subseteq \theta$ . Then  $T = \sum_{\theta} T_{\theta}$  has its Fourier-Wigner transform supported on Ω. Thus, by the corollary

$$\left\|\sum_{\theta\in\mathcal{P}_\Omega} T_\theta\right\|_{\mathcal{S}^p} \leq C(\Omega) \left\|\sum_{\theta\in\mathcal{P}_\Omega} T_\theta\star(g\otimes h)\right\|_{L^p(\mathbb{R}^{2d})} \leq C(\Omega)\mathcal{D}^{\mathscr{C}}_{p,q}\left(\sum_{\theta\in\mathcal{P}_\Omega} \|T_\theta\star(g\otimes h)\|^q_{L^p(\mathbb{R}^{2d})}\right)^{\frac{1}{q}},$$

as the Function  $F_{\theta} = T_{\theta} \star (g \otimes h)$  has its symplectic Fourier transform supported on  $\theta$ .

Applying Werner-Young's convolution theorem gives

$$\left\|\sum_{\theta\in\mathcal{P}_{\Omega}}\mathcal{T}_{\theta}\right\|_{\mathcal{S}^{p}}\leq C(\Omega)\mathcal{D}_{p,q}^{\mathscr{C}}\left(\sum_{\theta\in\mathcal{P}_{\Omega}}\|\mathcal{T}_{\theta}\star(g\otimes h)\|_{L^{p}(\mathbb{R}^{2d})}^{q}\right)^{\frac{1}{q}}\leq C(\Omega)\mathcal{D}_{p,q}^{\mathscr{C}}\left(\sum_{\theta\in\mathcal{P}_{\Omega}}\|\mathcal{T}_{\theta}\|_{\mathcal{S}^{p}}^{q}\right)^{\frac{1}{q}},$$

as g, h are  $L^2$ -normalised.

# Thank you for your attention

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